

**Fourier Series and
Transforms**

Revision Lecture

- The Basic Idea
- Real v Complex
- Series v Transform
- Fourier Analysis
- Power Conservation
- Gibbs Phenomenon
- Coefficient Decay Rate
- Periodic Extension
- Dirac Delta Function
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- Convolution
- Correlation

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Periodic signals can be written as a sum of sine and cosine waves:

$$u(t) = \frac{a_0}{2} + \sum_{n=1}^{\infty} (a_n \cos 2\pi n F t + b_n \sin 2\pi n F t)$$

The Basic Idea

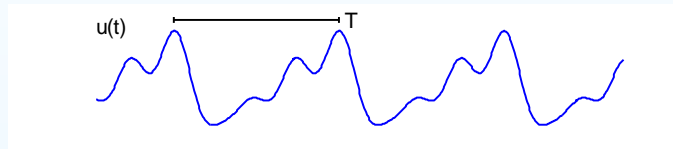
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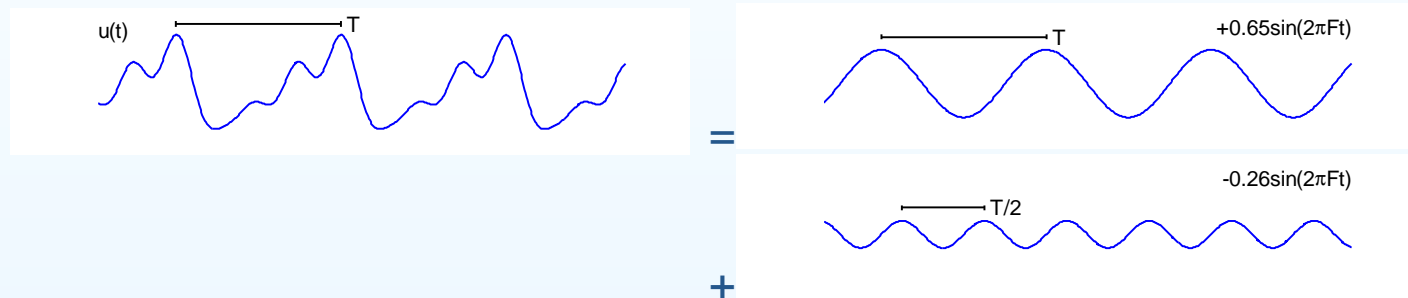
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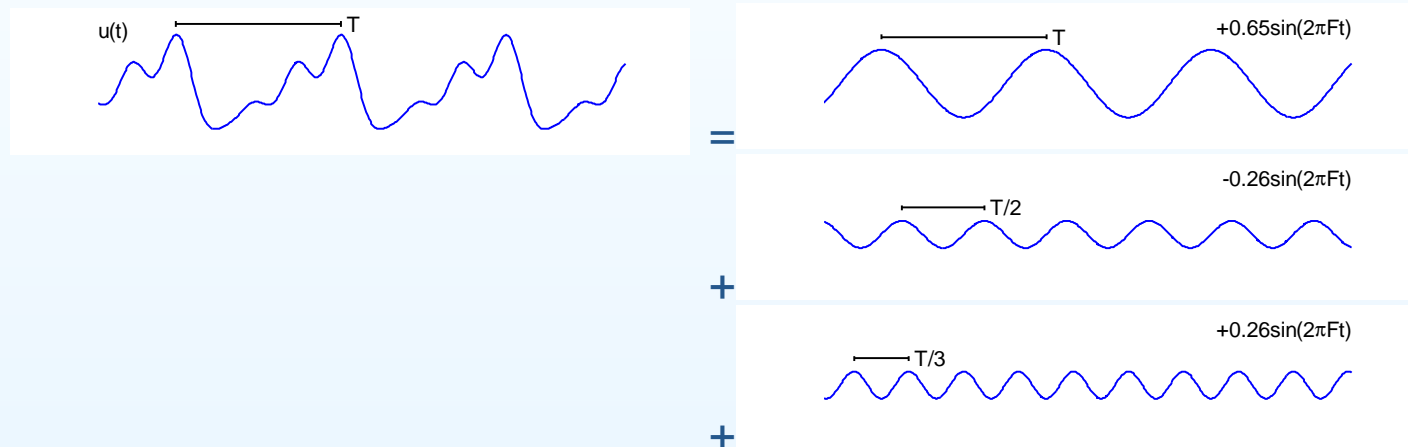
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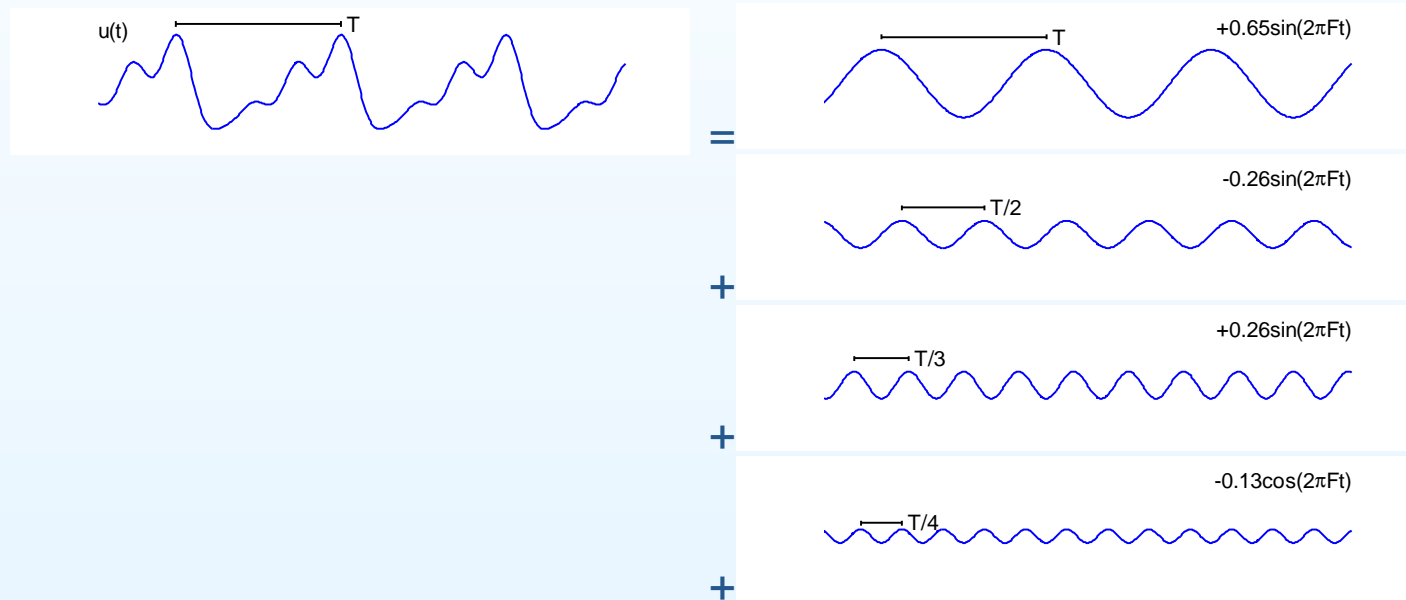
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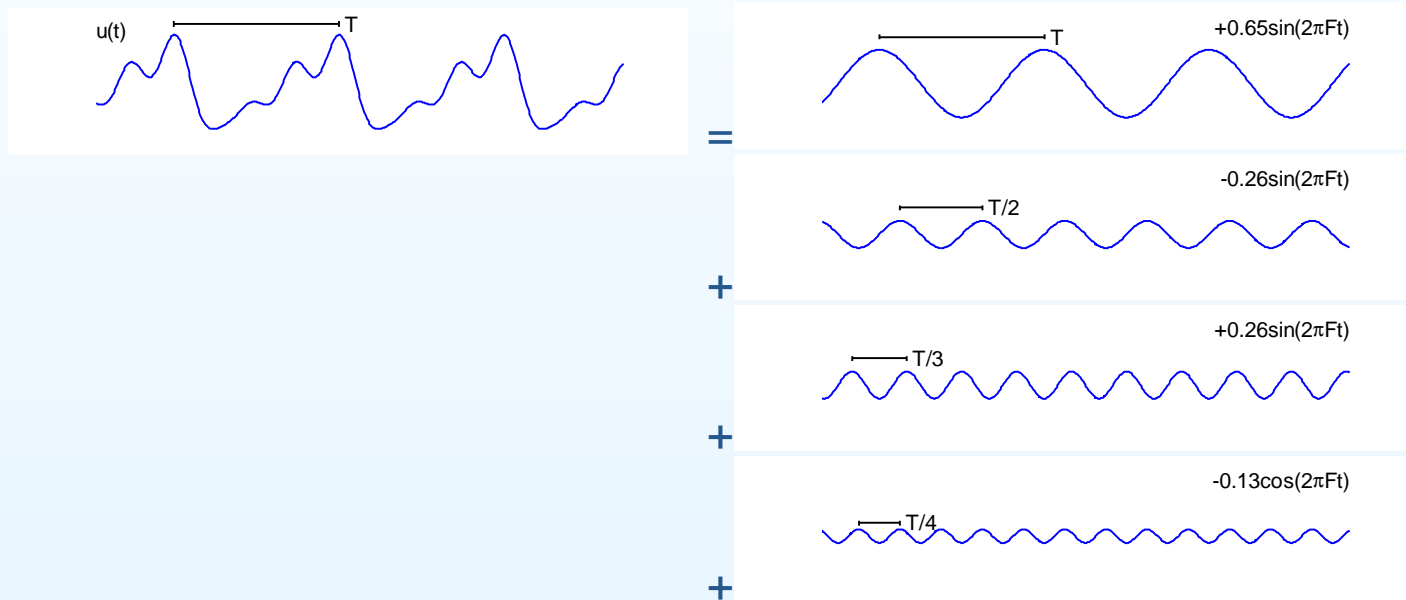
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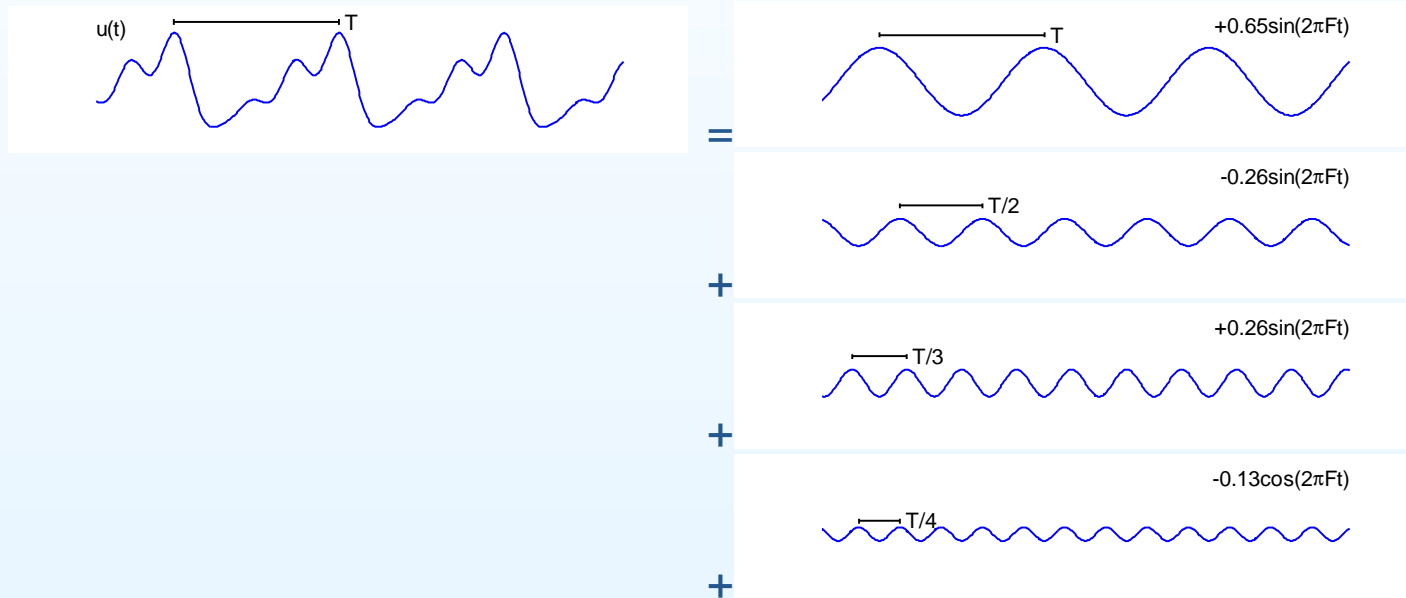
Fundamental Period: the smallest $T > 0$ for which $u(t + T) = u(t)$.

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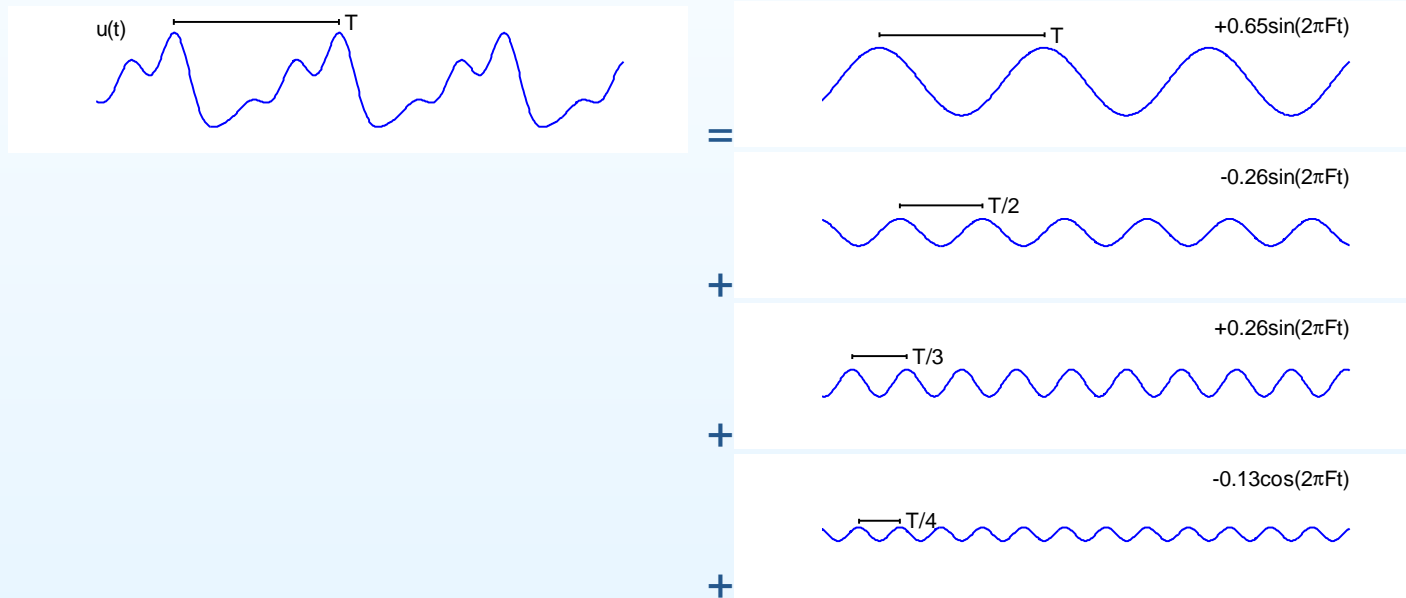
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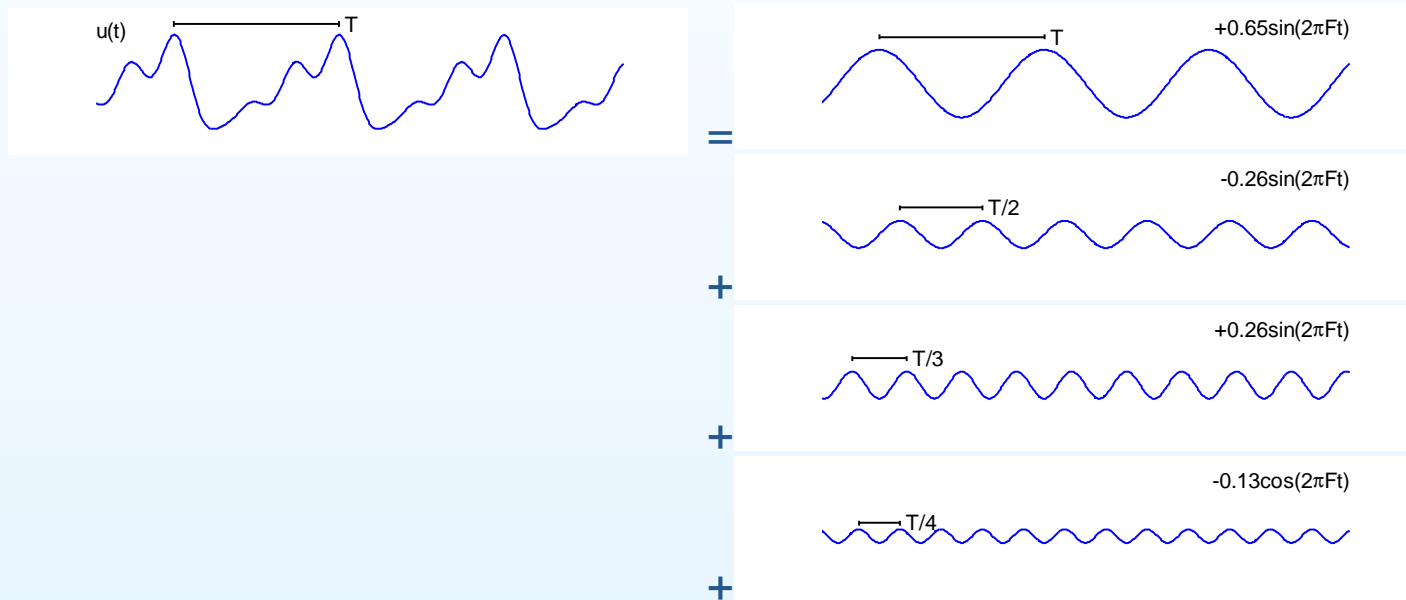
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Fundamental Period: the smallest $T > 0$ for which $u(t + T) = u(t)$.
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Some waveforms need infinitely many harmonics (countable infinity).

Real versus Complex Fourier Series

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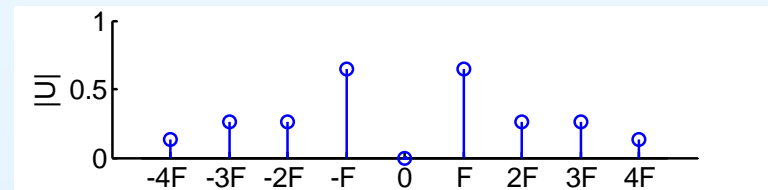
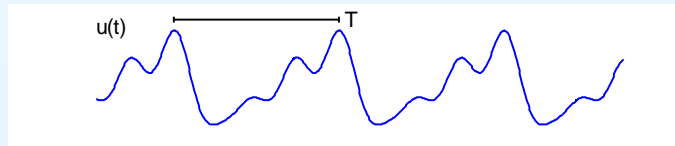
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Plot the **magnitude** spectrum

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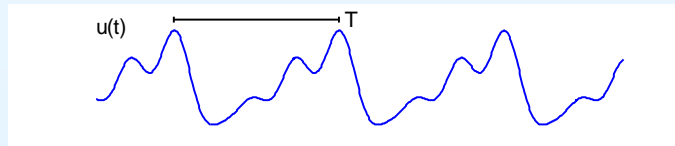
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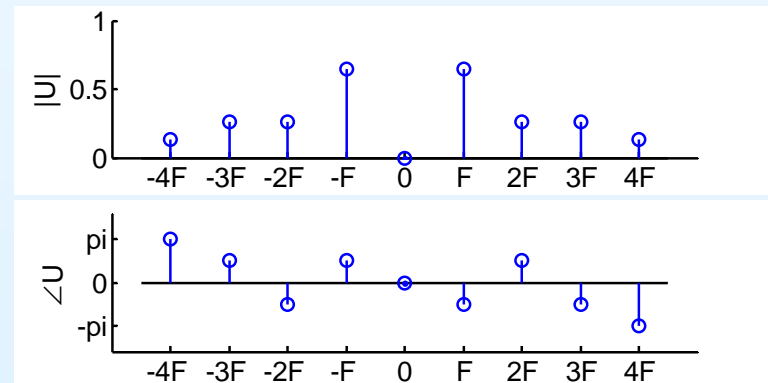
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Plot the **magnitude** spectrum
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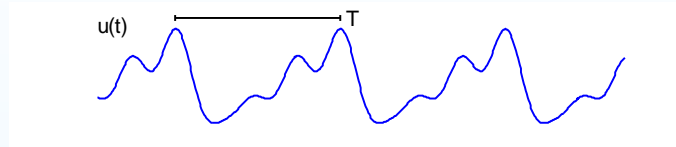
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- **Periodic** signals



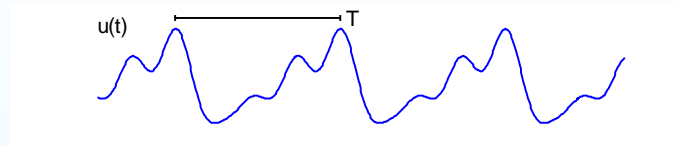
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- **Periodic** signals \rightarrow **Fourier Series**



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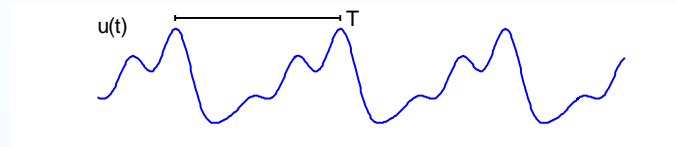
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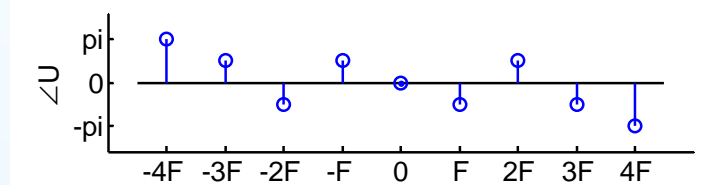
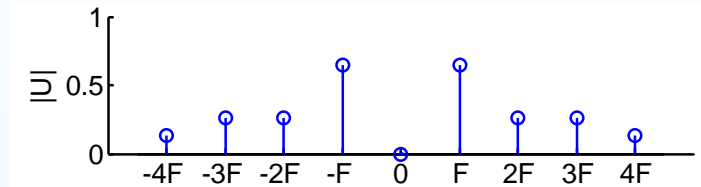
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- **Periodic** signals \rightarrow **Fourier Series** \rightarrow **Discrete** spectrum



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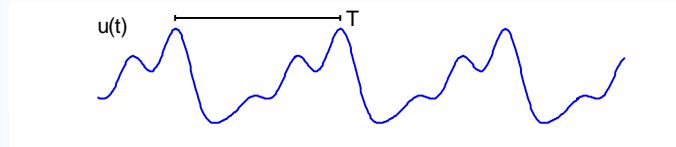
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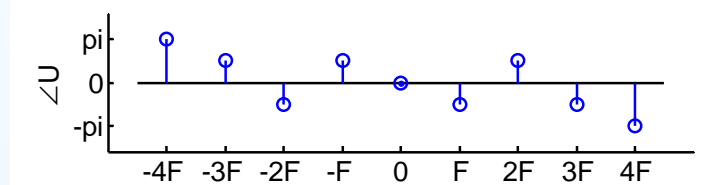
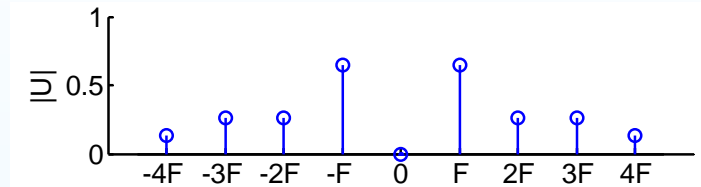
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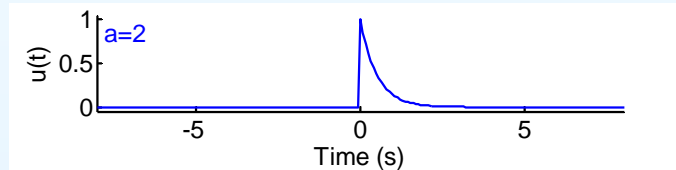
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- **Aperiodic** signals



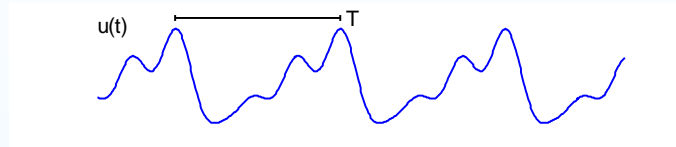
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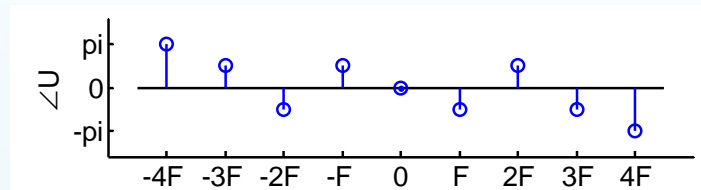
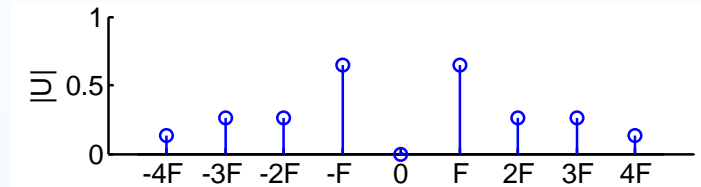
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- Dirac Delta Function
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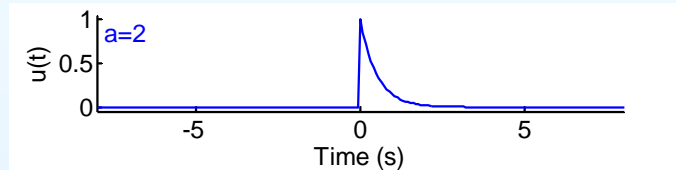
- **Periodic signals** → **Fourier Series** → **Discrete spectrum**



$$u(t) = \sum_{n=-\infty}^{\infty} U_n e^{i2\pi n F t}$$



- **Aperiodic signals** → **Fourier Transform**



$$u(t) = \int_{f=-\infty}^{\infty} U(f) e^{i2\pi f t} df$$

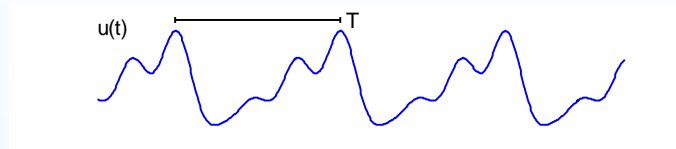
Fourier Series versus Fourier Transform

Fourier Series and Transforms

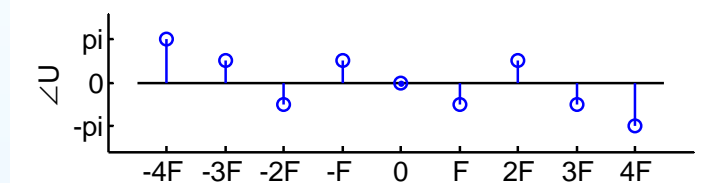
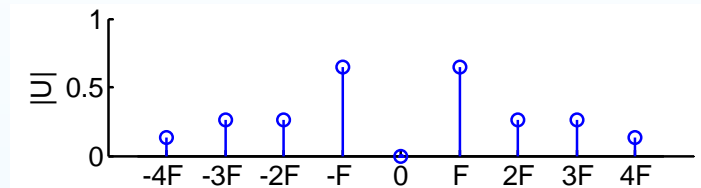
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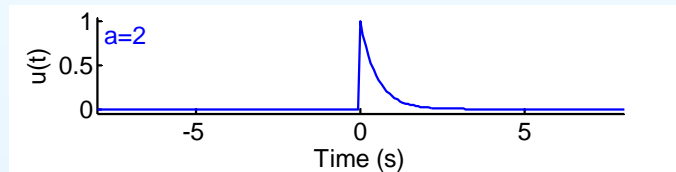
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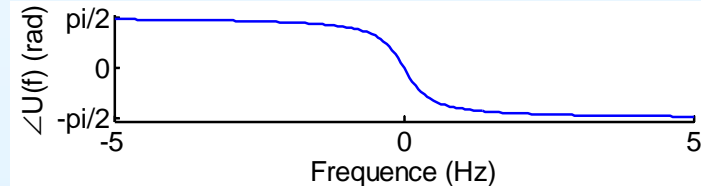
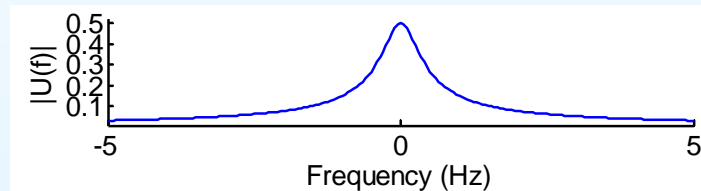
$$u(t) = \sum_{n=-\infty}^{\infty} U_n e^{i2\pi n F t}$$



- **Aperiodic signals** → **Fourier Transform** → **Continuous Spectrum**



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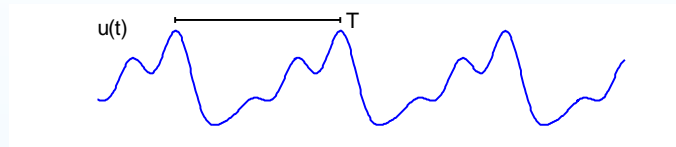
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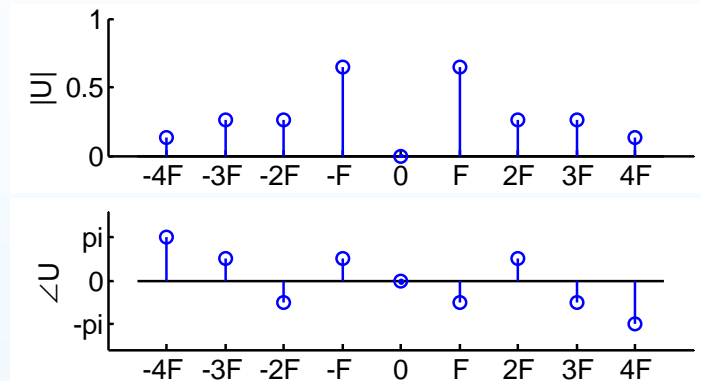
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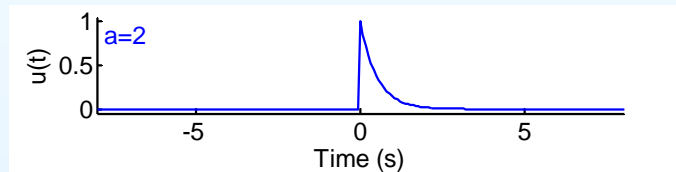
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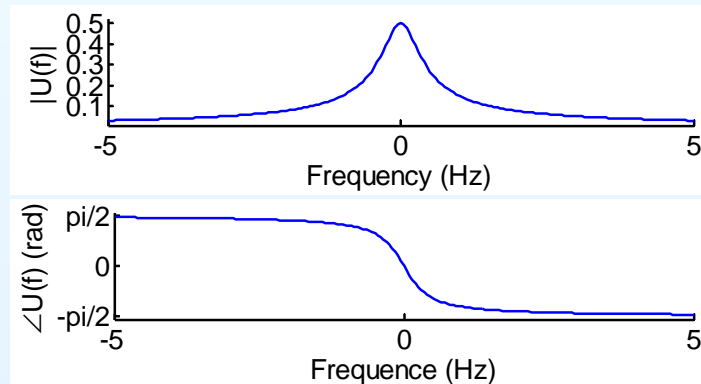
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- **Aperiodic signals** → **Fourier Transform** → **Continuous Spectrum**



$$u(t) = \int_{f=-\infty}^{\infty} U(f) e^{i2\pi f t} df$$



- Both types of spectrum are **conjugate symmetric**.
- If $u(t)$ is periodic, its Fourier transform consists of Dirac δ functions with amplitudes $\{U_n\}$.

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Fourier Series: $u(t) = \sum_{n=-\infty}^{\infty} U_n e^{i2\pi n F t}$

Fourier Analysis = “how do you work out the Fourier coefficients, U_n ?”

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$$\text{Key idea: } \langle e^{i\omega t} \rangle = \langle \cos \omega t + i \sin \omega t \rangle = \begin{cases} 1 & \text{if } \omega = 0 \\ 0 & \text{otherwise} \end{cases}$$

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So, to find a particular coefficient, U_m , we work out

$$\langle u(t) e^{-i2\pi m F t} \rangle = \langle \left(\sum_{n=-\infty}^{\infty} U_n e^{i2\pi n F t} \right) e^{-i2\pi m F t} \rangle$$

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Calculate the average by integrating over any integer number of periods

$$U_m = \langle u(t) e^{-i2\pi m F t} \rangle = \frac{1}{T} \int_{t=0}^T u(t) e^{-i2\pi m F t} dt$$

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Notice the **negative sign** in Fourier analysis: in order to extract the term in the series containing $e^{+i2\pi m F t}$ we need to multiply by $e^{-i2\pi m F t}$.

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$$\langle |U_n e^{i2\pi n F t}|^2 \rangle = \langle |U_n|^2 |e^{i2\pi n F t}|^2 \rangle = |U_n|^2$$

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$$\text{Truncated Fourier Series: } u_N(t) = \sum_{n=-N}^N U_n e^{i2\pi n F t}$$

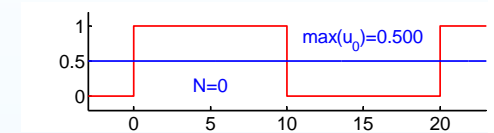
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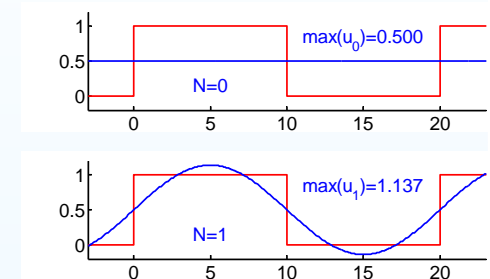
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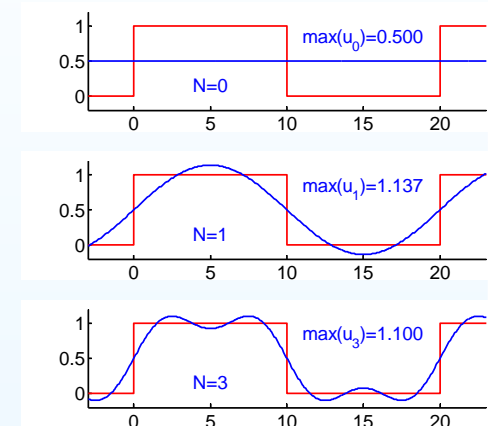
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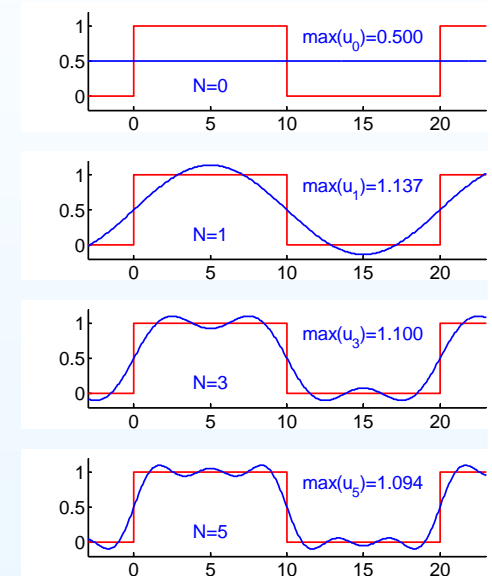
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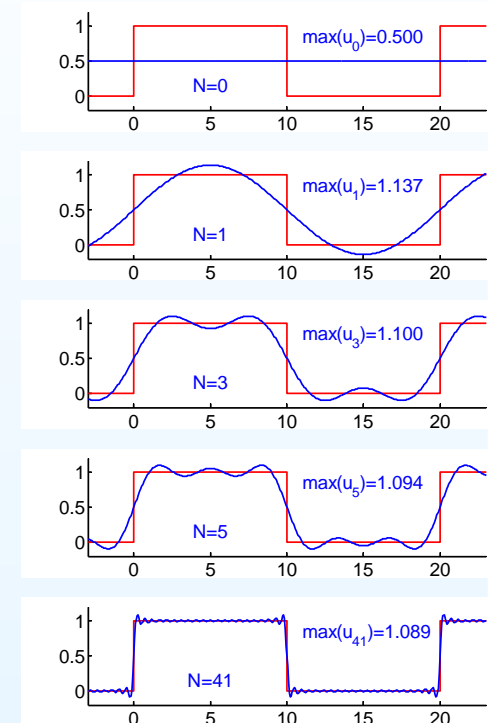
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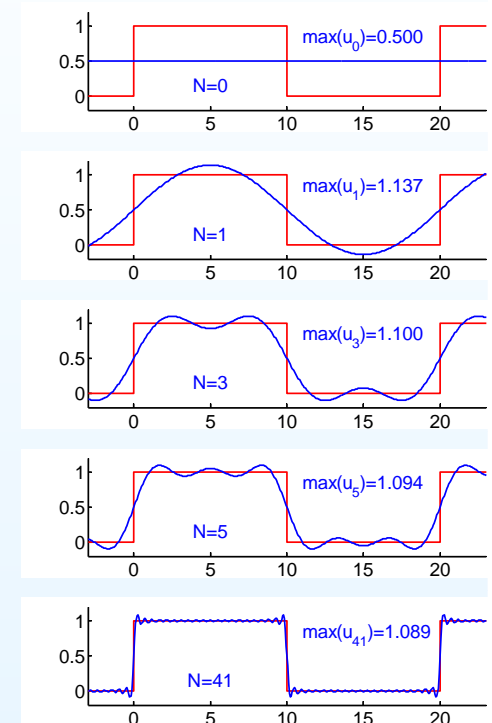
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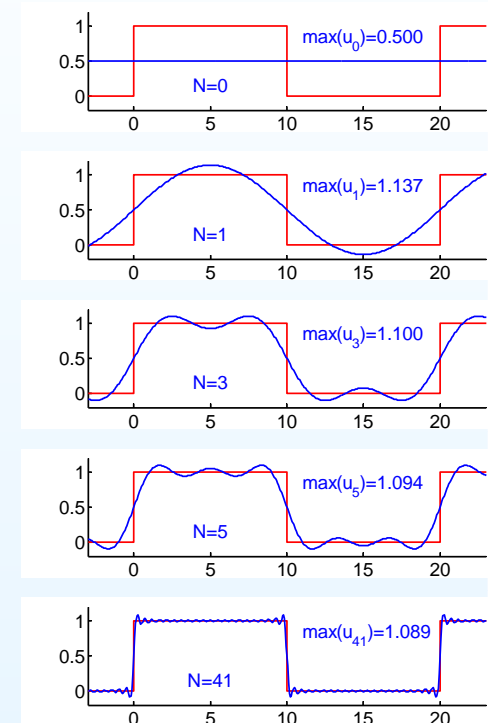
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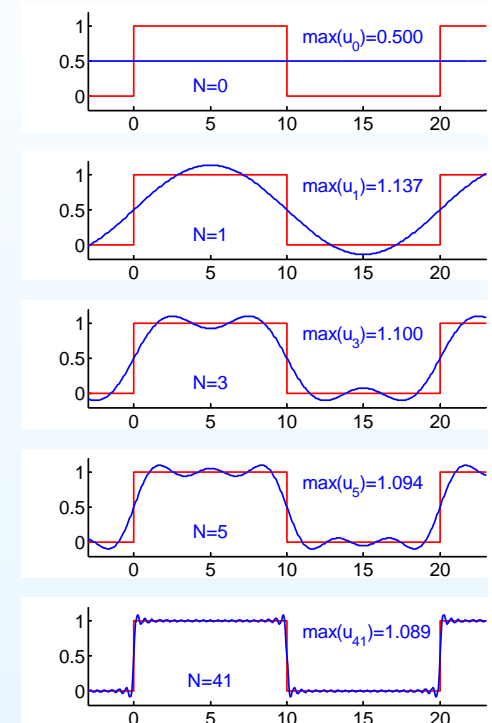
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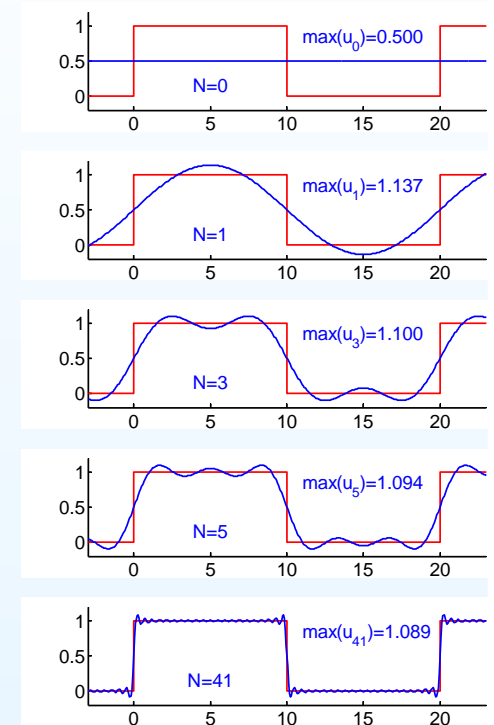
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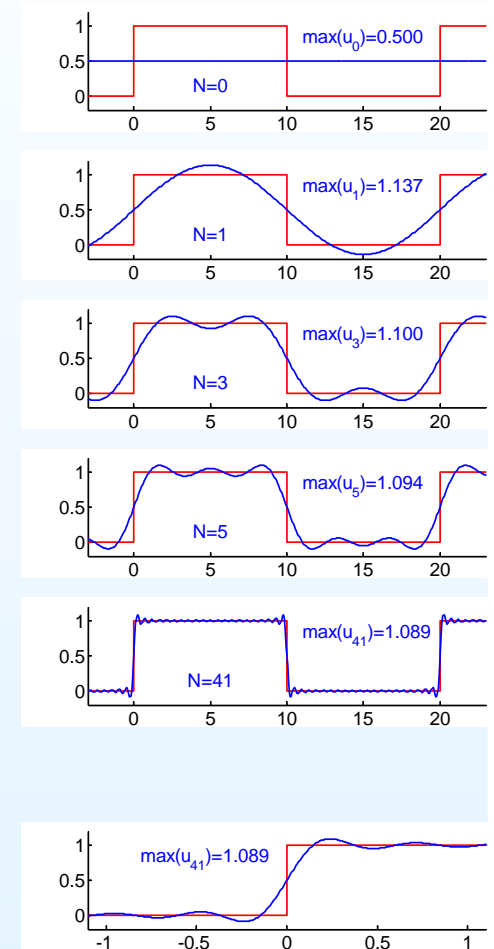
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If $u(t_0)$ has a discontinuity of height h then:

- $u_N(t_0) \rightarrow$ **the midpoint** of the discontinuity as $N \rightarrow \infty$.



[Enlarged View: $u_{41}(t)$]

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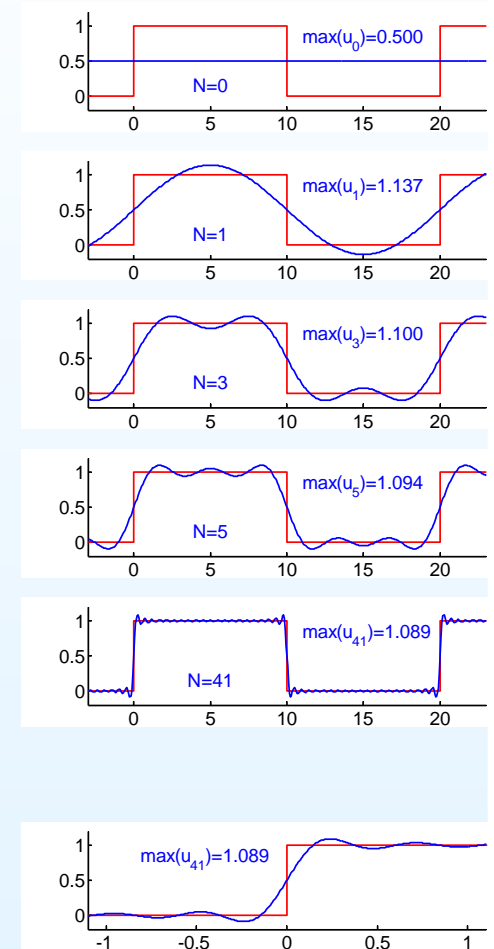
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If $u(t_0)$ has a discontinuity of height h then:

- $u_N(t_0) \rightarrow$ **the midpoint** of the discontinuity as $N \rightarrow \infty$.
- $u_N(t)$ **overshoots** by $\approx \pm 9\% \times h$ at $t \approx t_0 \pm \frac{T}{2N+1}$.



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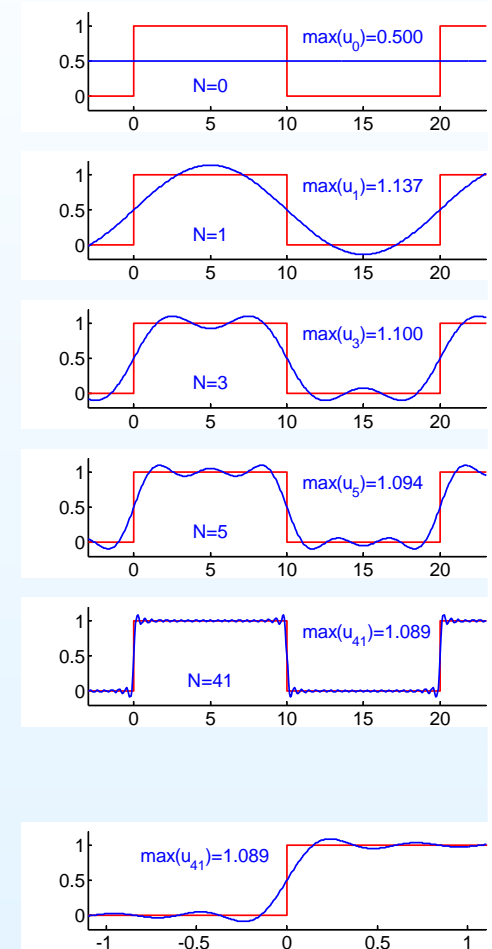
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- For large N , the overshoots move closer to the discontinuity but **do not decrease in size**.



[Enlarged View: $u_{41}(t)$]

Coefficient Decay Rate

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$$v(t) = \int_0^t u(\tau) d\tau \Rightarrow V_n = \frac{1}{i2\pi n F} U_n$$

provided $U_0 = V_0 = 0$.

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$$w(t) = \frac{du(t)}{dt} \Rightarrow W_n = i2\pi n F \times U_n$$

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If the coefficients, U_n , decrease rapidly then only a few terms are needed for a good approximation.

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Periodic Extension

If $u(t)$ is only defined over a finite range, $[0, B]$, we can make it periodic by defining $u(t \pm B) = u(t)$.

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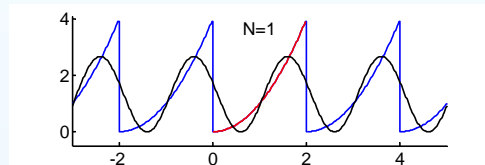
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Example: $u(t) = t^2$ for $0 \leq t < 2$



Periodic Extension

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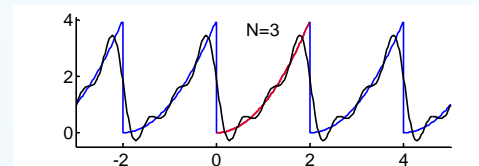
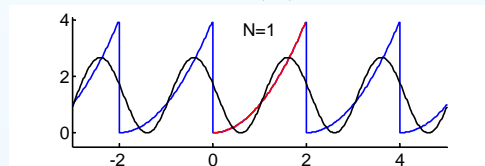
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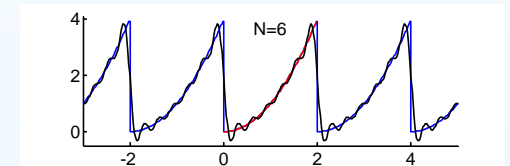
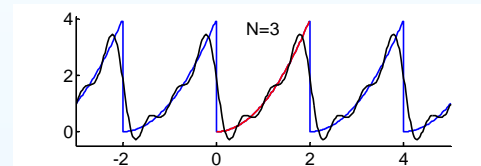
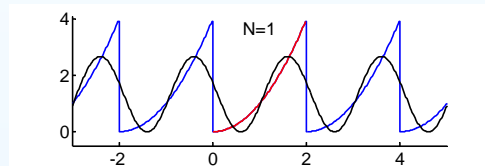
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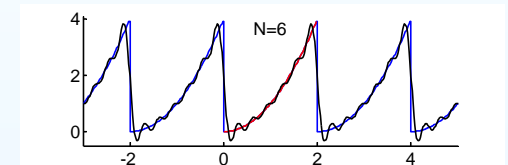
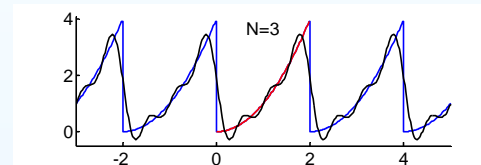
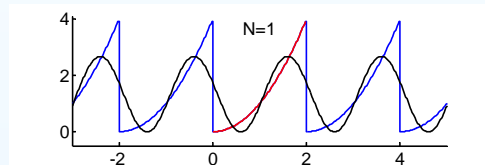
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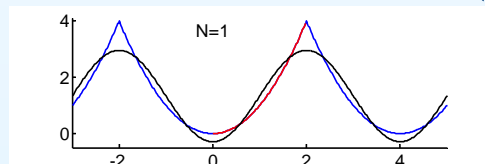
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Symmetric extension:

- To avoid a discontinuity at $t = T$, we can instead make the period $2B$ and define $u(-t) = u(+t)$.



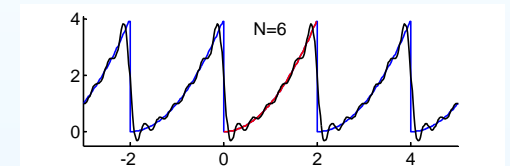
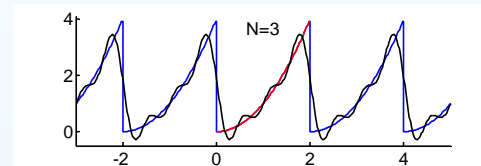
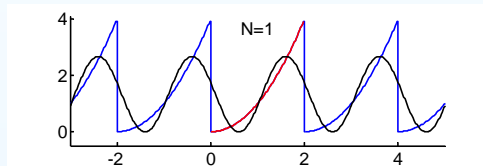
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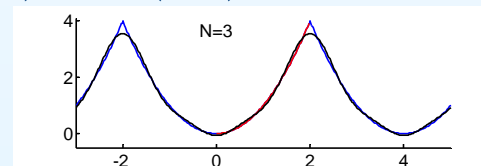
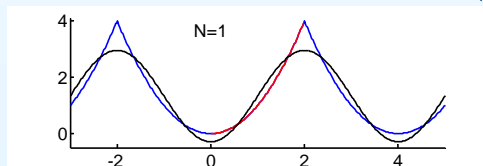
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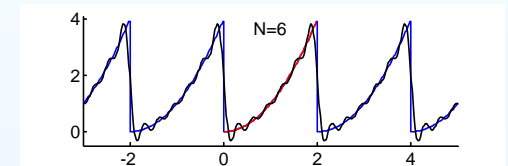
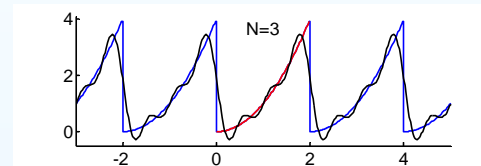
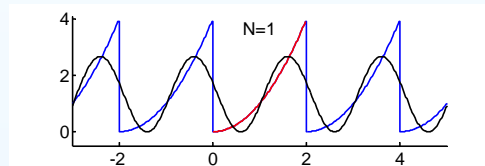
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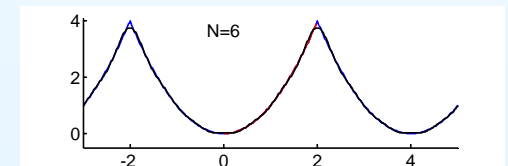
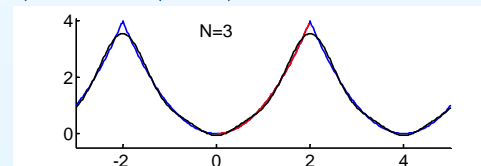
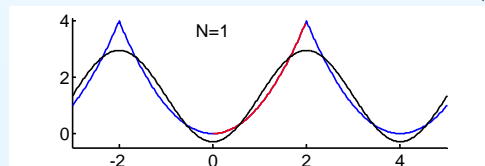
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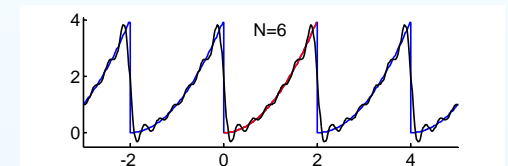
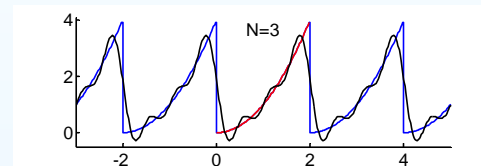
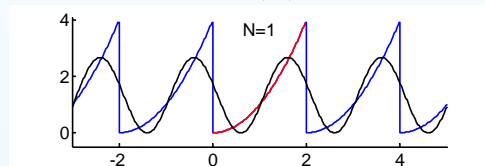
Periodic Extension

- The Basic Idea
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- Gibbs Phenomenon
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If $u(t)$ is only defined over a finite range, $[0, B]$, we can make it periodic by defining $u(t \pm B) = u(t)$.

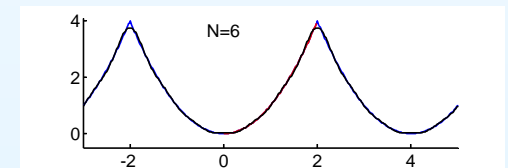
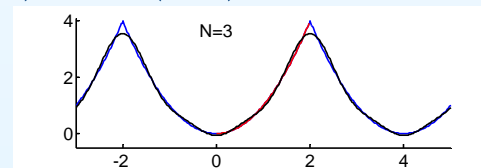
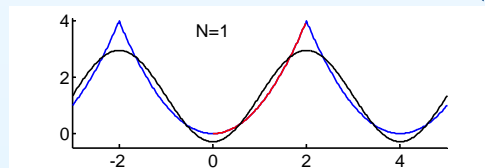
- Coefficients are given by $U_n = \frac{1}{B} \int_0^B u(t) e^{-i2\pi n F t} dt$.

Example: $u(t) = t^2$ for $0 \leq t < 2$



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- To avoid a discontinuity at $t = T$, we can instead make the period $2B$ and define $u(-t) = u(+t)$.



- Symmetry around $t = 0$ means coefficients are **real-valued** and **symmetric** ($U_{-n} = U_n^* = U_n$).

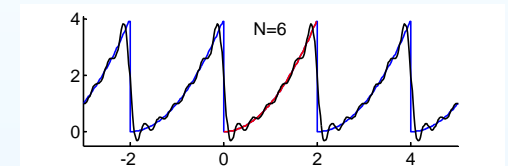
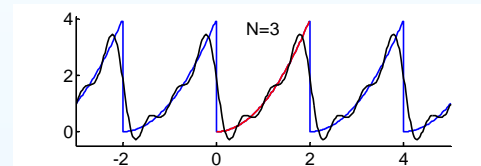
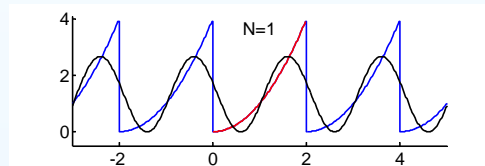
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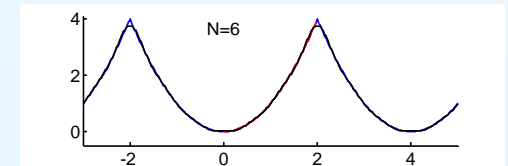
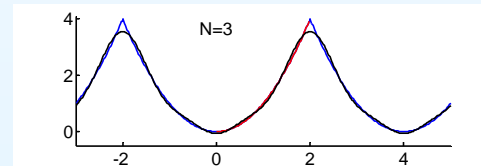
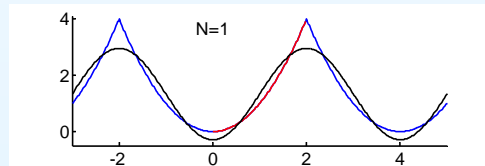
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- To avoid a discontinuity at $t = T$, we can instead make the period $2B$ and define $u(-t) = u(+t)$.



- Symmetry around $t = 0$ means coefficients are **real-valued** and **symmetric** ($U_{-n} = U_n^* = U_n$).
- Still have a first-derivative discontinuity at $t = B$ but now we have **no Gibbs phenomenon** and coefficients $\propto n^{-2}$ instead of $\propto n^{-1}$ so approximation error power decreases more quickly.

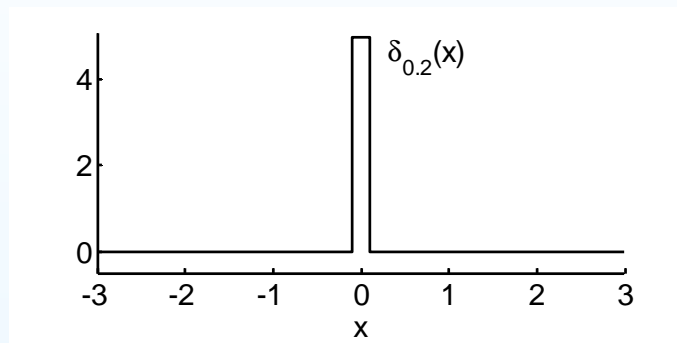
Dirac Delta Function

Fourier Series and
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$\delta(x)$ is the limiting case as $w \rightarrow 0$ of a pulse w wide and $\frac{1}{w}$ high



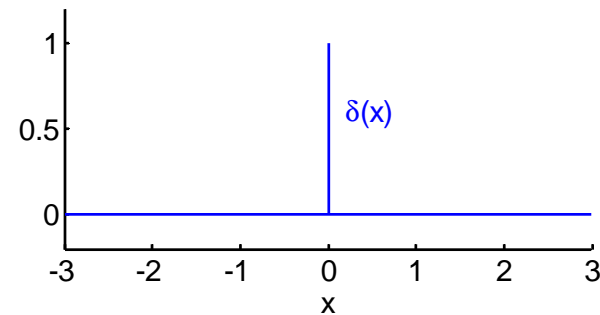
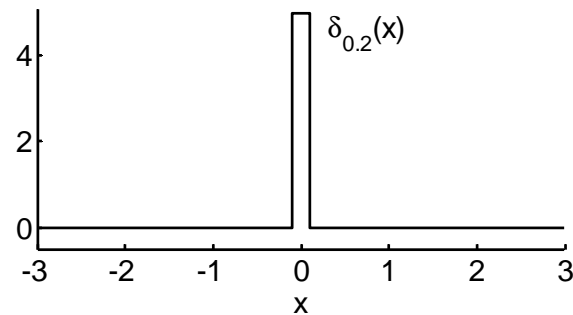
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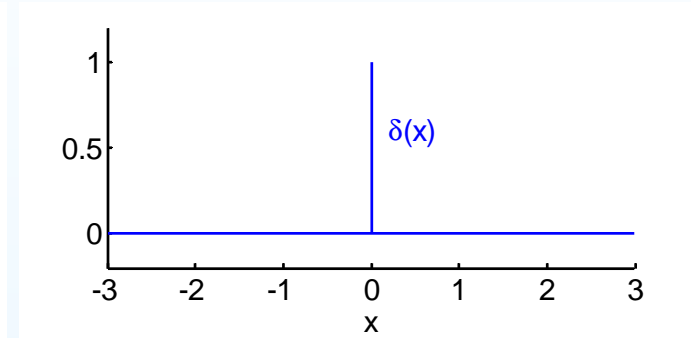
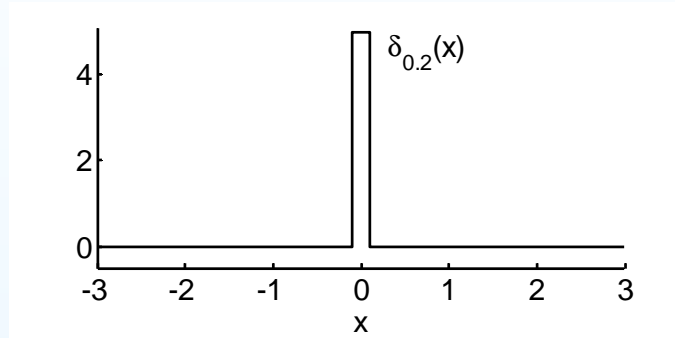
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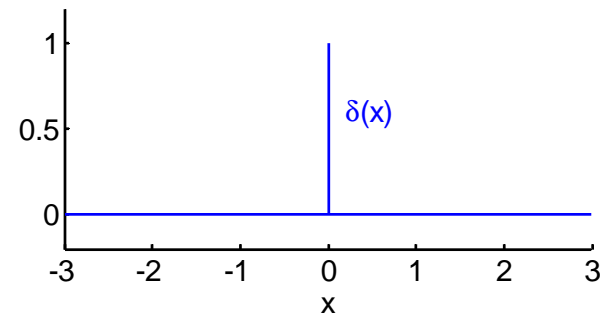
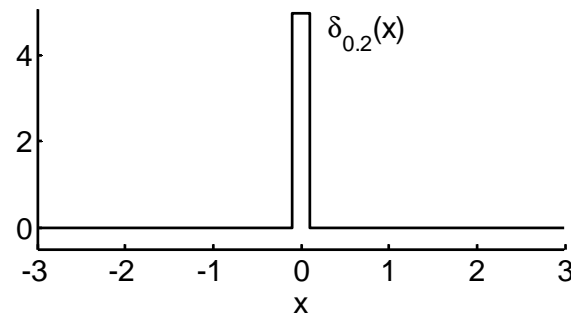
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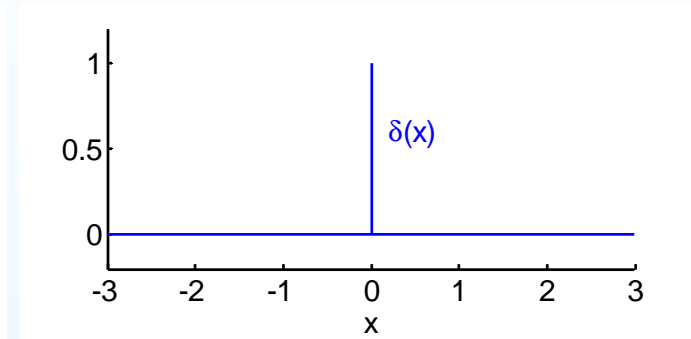
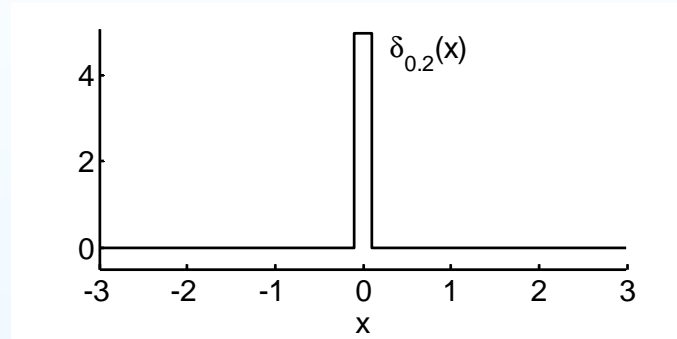
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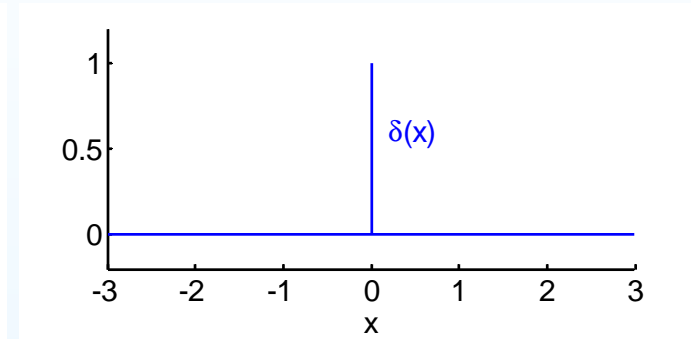
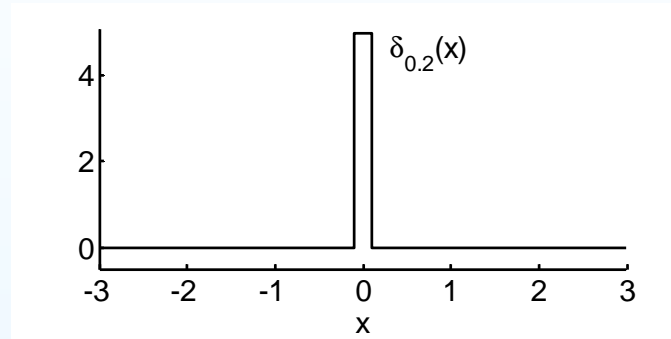
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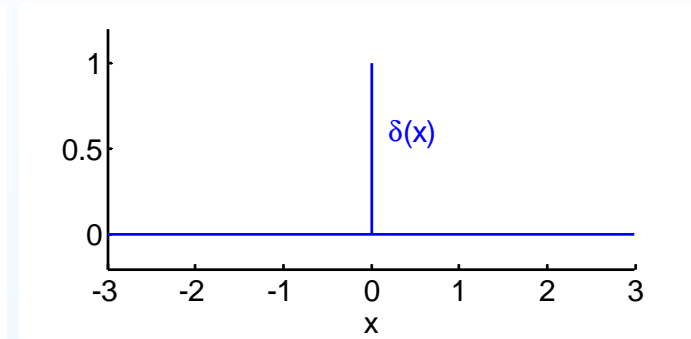
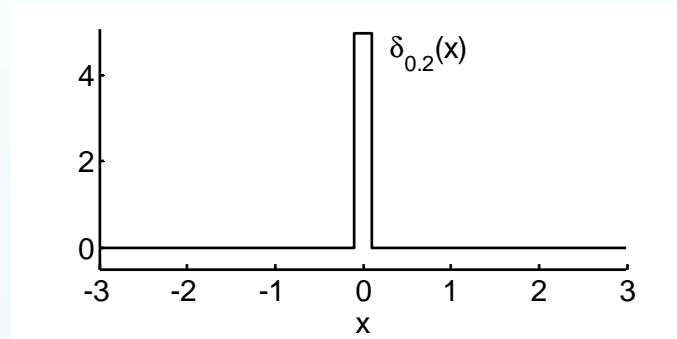
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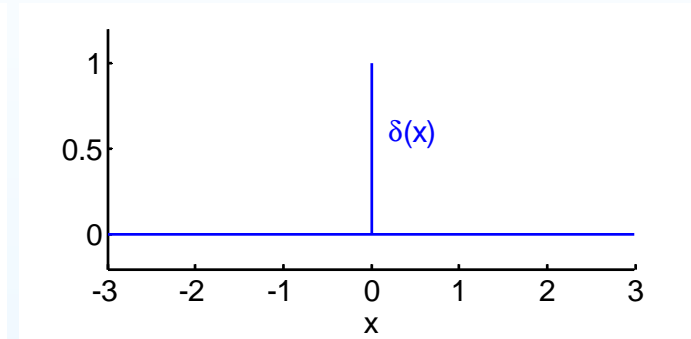
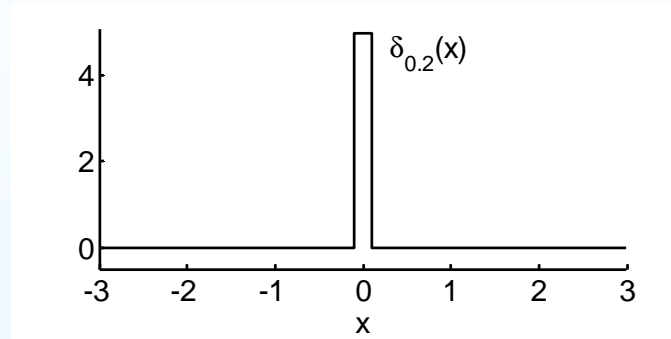
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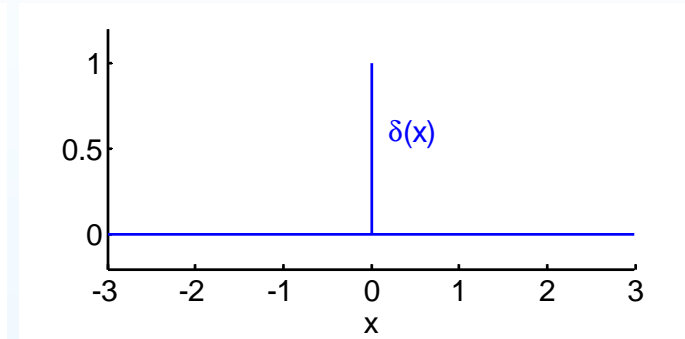
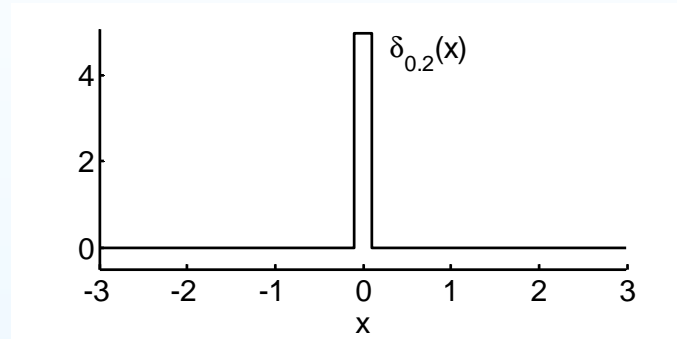
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- We plot $h\delta(x)$ as a pulse of height $|h|$ (instead of its true height of ∞)

Fourier Transform

Fourier Series and
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$$u(t) = \int_{-\infty}^{\infty} U(f)e^{i2\pi ft} df$$
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Fourier Transform

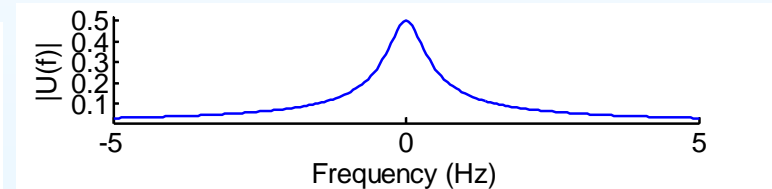
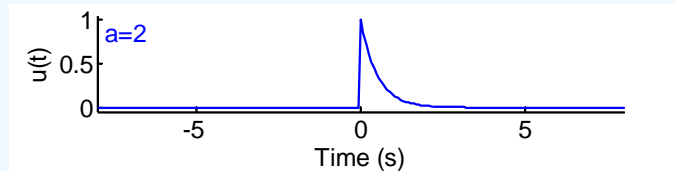
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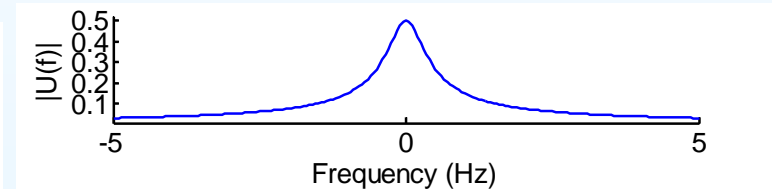
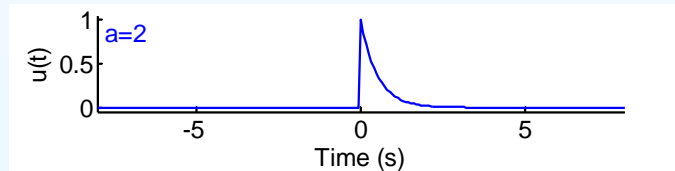
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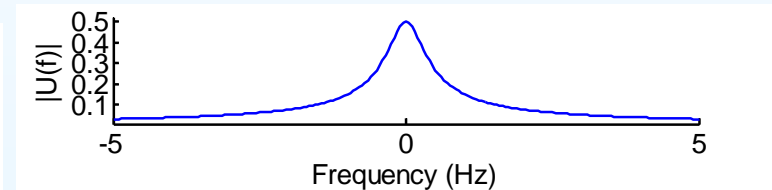
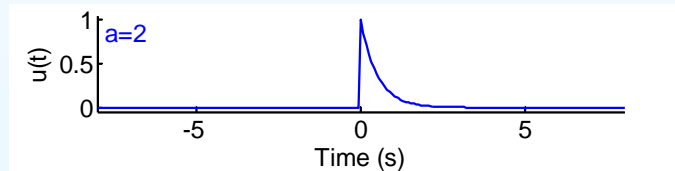
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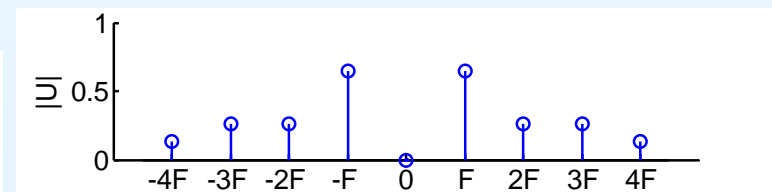
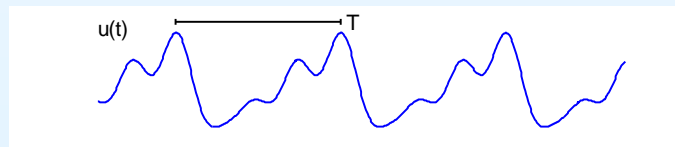
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- **Periodic Signals** \rightarrow Dirac δ functions at harmonics.
Same complex-valued amplitudes as U_n from Fourier Series



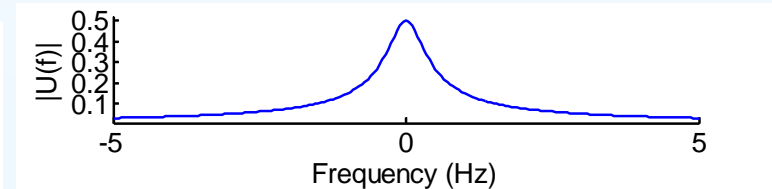
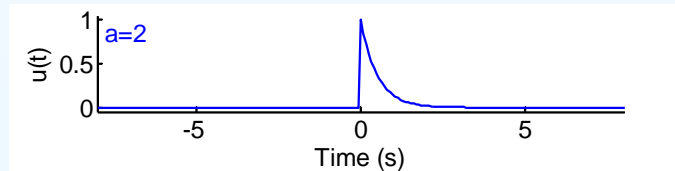
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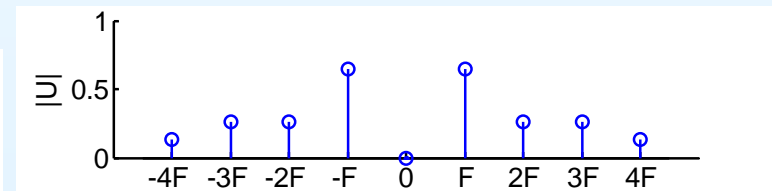
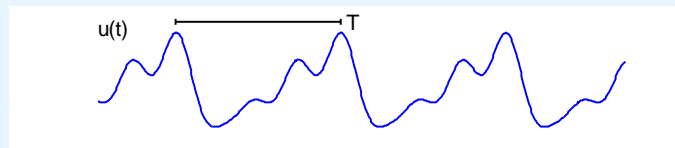
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- $E_u = \infty$ but ave power is $P_u = \langle |u(t)|^2 \rangle = \sum_{n=-\infty}^{\infty} |U_n|^2$

Convolution

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Multiplication in either the time or frequency domain
is equivalent to **convolution** in the other domain:

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Convolution

Fourier Series and
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Revision Lecture

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- Real v Complex
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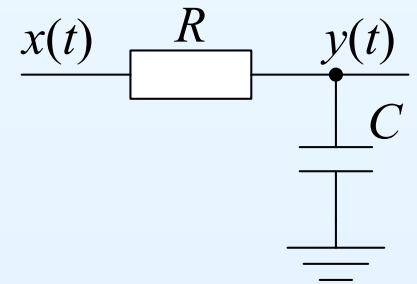
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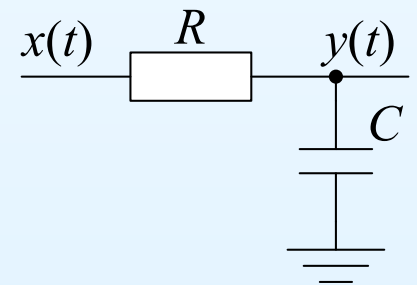
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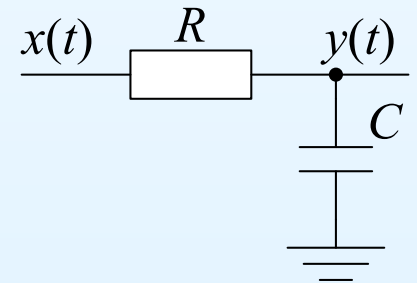
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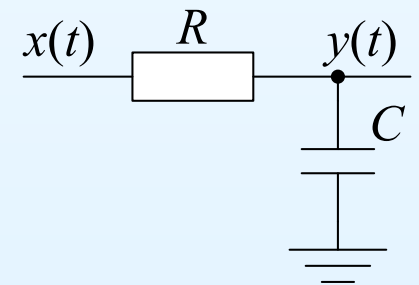
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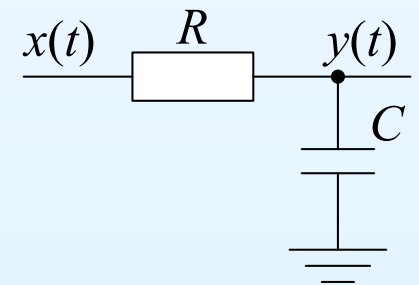
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- For all values of t : $|w(t)|^2 \leq E_u E_v$
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- **Cross-correlation** is used to find the time shift, t_0 , at which two signals match and also how well they match.
- **Auto-correlation** is the cross-correlation of a signal with itself: used to find the period of a signal (i.e. the time shift where it matches itself).